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A Fast Radio Burst Backend for the Tianlai Dish Pathfinder Array

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Abstract

The Tianlai Dish Pathfinder Array is a radio interferometer array consisting of 16 six-meter dish antennas. The original digital backend integration time is at the seconds level, designed for H I intensity mapping experiment. A new digital backend with millisecond response is added to enable it to search for fast radio burst during its observations. The design and calibration of this backend, and the real time search pipeline for it are described in this paper. It is capable of forming 16 digital beams for each linear polarization, covering an area of 19.6 square degrees. The search pipeline is capable of searching for, recording and classifying FRBs automatically in real time. In commissioning, we succeeded in capturing the signal pulses from the pulsars PSR B0329+54 and B2021+51.

Key words: instrumentation: miscellaneous – techniques: interferometric – telescopes – radio continuum: general – (stars:) pulsars: general

1. Introduction

Since their discovery (Lorimer et al. 2007; Thornton et al. 2013), fast radio bursts (FRBs) have been an intensely studied topic in radio astronomy. The cosmological origins of them are established after FRB 121102 was localized to a host galaxy with z = 0.193 (Bassa et al. 2017; Chatterjee et al. 2017; Marcote et al. 2017). The astrophysical origin of FRBs and radiation mechanism are still unsolved problems, but in the past few years significant progresses have been made, with many more events observed by the various telescopes. Up to now, over 800 fast radio bursts, including 25 repeaters, have been publicly announced.¹¹ The discovery of FRB 20200428 (CHIME/FRB Collaboration et al. 2020), which is originated from the known galactic magnetar SGR J1935+2154, indicates that magnetars are probable sources for at least some FRBs (e.g., Wang et al. 2020).

The FRB flash typically lasts a few milliseconds, so that the radio telescope must have a sub-millisecond sampling rate to detect the FRBs. Given the large amount of data received by the interferometer array, it is impractical to store all of the data at such high time resolution. Instead, the data within a short time frame are searched in real time for a possible FRB, and the high time resolution data are stored for further analysis if a candidate is found, or else discarded if no FRB is found.

This paper describes the FRB digital backend of the Tianlai Dish Pathfinder Array. The Tianlai experiment is dedicated to test the key technologies for 21 cm intensity mapping method (Chen 2012; Xu et al. 2014), which consists a cylinder pathfinder array (Zhang et al. 2016; Li et al. 2020, 2021; Sun et al. 2022) and a dish pathfinder array (Zhang et al. 2016; Das et al. 2018; Wu et al. 2021; Phan et al. 2022). It consists of 16 onaxis dish antennas of 6-meter aperture, equipped with dual linear polarization receivers, its basic parameters are listed in Table 1, and the array configuration is shown in Figure 1. The original Tianlai correlator is only designed to output visibilities with a ~ 1 second time sampling, which is then processed offline for sky imaging (Zuo et al. 2021). However, with a field of view (FoV) of \sim 19.6 square degrees, the array is potentially capable of discovering large numbers of FRBs while doing its H I survey (Perdereau et al. 2022), if it is equipped with digital backend which has millisecond time resolution to perform FRB searches.

Below, we first present the operation principle and design of the Tianlai Dish Array FRB backend, then describe our commissioning experiment and test results, followed by an estimate of its sensitivity and the expected number of FRBs. Finally we

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¹¹ A catalog is available at https://www.herta-experiment.org/frbstats/.



Figure 1. Tianlai dish antenna layout with antenna number and ground coordinates with respect to a reference point (unit: meter).

summarize the results. For the antennas in the Tianlai dish array, we use labels explained in Wu et al. (2021) and also shown in Figure 1.

2. System Design

The FRBs occur randomly in any direction of the sky, so a blind FRB survey can be conducted by pointing the telescope in any direction and searching for burst events in the output. The FRB signal is dispersed by the interstellar medium, so that the higher frequency part arrives slightly earlier, with the time delay (in units of seconds) with respect to $\nu \rightarrow \infty$ given by

$$\Delta t \approx 4.148 \times 10^{-3} \frac{\text{DM}}{\text{pc cm}^{-3}} \left(\frac{\nu}{\text{GHz}}\right)^{-2}$$
(1)

where DM is the dispersion measure along the propagation path. The dispersion phenomenon can be used as a character to distinguish the FRB or pulsar signal from the radio frequency interference (RFI), which occurs much more frequently.

The conventional correlator for the Tianlai array produces interferometric visibilities, i.e., the short time integration of the cross-correlation of voltages of different receiver pairs, and the sky image is synthesized from such visibilities, which corresponds approximately to Fourier components of the sky radiation intensity (Thompson et al. 2017). However, there are a large number of visibilities (528 including auto-correlations), the individual sensitivity to a spatially localized source such as an FRB is not high. To facilitate the search of an FRB, we

 Table 1

 Basic Properties of Tianlai Dish Pathfinder Array

Antenna mount	Alt-Az pedestal
Number of Antenna	16
Dish diameter	6 m
f/D	0.37
Average SEFD	14.15 kJy
Latitude	44°15 N
Longitude	91°80 E
Observing frequency	685–810MHz

adopt a beam-forming approach for the FRB search engine. We digitally form multiple beams sensitive to different sky directions, the output data is then de-dispersed with a range of dispersion measure (DM) values and searched for pulses above the noise. This is realized with an independent digital system, which shares the analog frontend hardware of the radio array with the original correlator, but operates independently. Below we describe this digital system.

2.1. Beam Forming Principle

The voltage output from antenna unit *a* is:

$$\begin{aligned} \mathcal{E}_{a} &= |g_{a}|e^{j\phi_{a}}\int e^{-2\pi j \boldsymbol{n}\cdot\boldsymbol{u}_{a}} \\ &\times A_{i}(\boldsymbol{n})I(\boldsymbol{n})d^{2}\boldsymbol{n} + \eta_{a}, \end{aligned} \tag{2}$$

where $A_a(\mathbf{n})$ describes the voltage response for antenna a, \mathbf{u}_a is the position vector of the antenna in units of wavelength, the integration is over sky directions, and $g_a = |g_a|e^{j\phi_a}$ is the complex gain of the instrument for unit a, η_a represents the noise for that channel. For the Tianlai array, the instrument phase ϕ_a mainly comes from the optical cable, though the feed and amplifier/bandpass filter may also contribute. We can form a digital beam from the outputs of the array by adding the voltages with different complex weights,

$$S(\mathbf{k}) = \sum_{a} w_a(\mathbf{k}) \mathcal{E}_a \tag{3}$$

where $w_a(\mathbf{k})$ denotes the complex weight of antenna *a* for beam direction \mathbf{k} , which we can set as

$$w_a(\mathbf{k}) = \frac{1}{|g_a|} e^{-j\phi_a} e^{2\pi j \mathbf{k} \cdot \mathbf{u}_a},\tag{4}$$

so that ideally,

$$S(\mathbf{k}) = \sum_{a} \int e^{-2\pi j (\mathbf{n} - \mathbf{k}) \cdot \mathbf{u}_{a}} \\ \times A_{i}(\mathbf{n}) I(\mathbf{n}) d^{2} \mathbf{n}.$$
(5)

In this integral, the phases of all antennas are equal for the direction n = k, so that their voltages add coherently, while for other directions the terms in the sum would be out of phase, so this output is most sensitive to the direction k. If a burst occurs in the direction q at some moment, and assume that sky in other



Figure 2. The profile for digitally formed beam in dB at 750 MHz. We assume all the primary beams are identical and pointed to the North Celestial Pole (NCP). Left: the digital beam centered at NCP. Right: the digital beam centered at the same elevation, but 2° East of NCP. The amplitudes of beams are normalized so that the beam centered at NCP has maximum amplitude of 0 dB.

directions remains nearly constant at this time, then the received signal is given by

$$S = S_0 + \Delta S \tag{6}$$

where

$$\Delta S = \sum_{a} e^{-2\pi j (\boldsymbol{q} - \boldsymbol{k}) \cdot \boldsymbol{u}_{a}} A_{i}(\boldsymbol{q}) I(\boldsymbol{q}), \qquad (7)$$

and S_0 is the S value without burst.

The primary beam of the Tianlai dish has been measured with unmanned air vehicle (Zhang et al. 2021). Here we simply model the primary beam as

$$A(\theta) = A_0 \operatorname{sinc}\left(\frac{\theta}{2\sigma}\right) \tag{8}$$

where θ is the angle between the direction of interest to the direction of antenna axis and $\sigma \approx 0^{\circ}.79$ corresponding to a full width at half maximum (FWHM) $\approx 4^{\circ}.4$. The digitally formed beam can point to any direction, but for different directions the beam shape would be somewhat different, as the antennas are fixed in position and the projected baselines are different for different directions.

In Figure 2 we show the synthesis beam centered at the North Celestial Pole (NCP) (Left), and the beam centered at a point at the same elevation of NCP but 2° to its East (Right). The FWHM of the main lobe of the synthesis beam is about 0°.6 for the east–west direction and 0°.8 for the north–south direction. There are also sidelobes which extend over the primary beam. Note that the profile of both the primary beam and the digitally formed beam varies with frequency, but for the current Tianlai arrays which have a fractional bandwidth of ~13% at 750 MHz the variation is moderate.

2.2. The System Setup

The FRB system shares the same analog hardware system with the original Tianlai dish correlator system, as shown in Figure 3. A detailed description of the analog system can be found in Wu et al. (2021), and here we give a brief summary. The Tianlai dish array consists of 16 steerable parabolic reflector antennas. Each antenna is equipped with a dual, linearpolarization feed, thus 32 polarized radio signals are collected in total. The signals are amplified by the low noise amplifiers (LNA), which are installed on the backside of the feed unit, and transmitted by a 15 m long cable down to the optic transmitters below the reflectors. In the Tianlai experiment, the station house where the signal chain and digital facilities are running are located in a village which is 6 km away from the antenna area. The radio signals are converted to optical signals at the antenna site and communicated through optical fibers to the station house, where they are converted back to radio frequency (RF). Then, the RF signals are bandpass filtered for a selected 100 MHz bandwidth (currently in the 700-800 MHz, but maybe shifted to other frequency bands). The signals are then down-converted to the intermediate frequency (IF) with a range of 135-235 MHz via a mixer. The IF signals are further amplified. Up to this point, the FRB system and the original correlator system share the same signal chain hardware.

Each of the IF signal is then split into two channels by a power splitter. The signal is amplified with an additional amplifier to its original magnitude to compensate for the power loss. One of signal is fed to the original digital correlator, while the other is fed to the FRB searching system. A schematic of the whole Tianlai Dish Array system is shown in Figure 3.

The hardware of the FRB digital backend consists of FPGAbased beamformers and GPU-based de-dispersion servers. The beamformer digitize the signal, and form multiple beams



Figure 3. Schematic of the Tianlai FRB-search and correlation system.

according to Equations (3)–(7). The output of the beamformer is send to the de-dispersion server, which de-disperse the signal with a range of DM values and check for possible FRB candidates. At present, the Tianlai Dish Array beamformer consists of two FPGA boards, each processes the input from one linear polarization of the feeds, and correspondingly there are also two GPU servers. The beam former time resolution is up to 98.304 μs with 1024 frequency channels and 122.07 kHz channel width. The FPGA-based beamformer and the real time processing of the beamformed data by the de-dispersion system are described in the next two subsections.

2.3. The FPGA Beamformer

As shown in Figure 3, the FPGA beamformer for the Tianlai dish array consists of two CASPER SNAP2¹² boards,¹³ each of them processes data from one of the linear polarizations. Each SNAP2 board has two analog-to-digital converter (ADC) boards, and each ADC converts analog input signal from eight units into digital data. Each SNAP2 board processes the raw data from both ADCs, and output the beamformed intensity data to a de-dispersing GPU server through four 10GbE network interface cards (NICs). For system testing, the ADCs are also capable of returning the real-time data directly, bypassing the data processing on the SNAP2 board.

The workflow of the beamformer is shown in Figure 4. The raw data collected by the ADCs needs to be calibrated for delay and phase compensation, due to the different cable lengths and the change of geometric delay when the antennas are pointing to different directions. For the Tianlai Dish Pathfinder Array, the maximum signal delay between two feeds is about 50 ns. This delay is compensated in two steps. The first step deals with the major part of the delay up to multiples of the sampling period which is 4ns. This is done by shifting the data block. Then, in the second step, after polyphase filtering and fast Fourier transforming the data to frequency domain, the residual of the delay is compensated by a phase factor, which is given by the sum of geometric delay which is computed for each pointing direction, and the instrument delay which is determined by calibration with astronomical source. Compensation coefficients are then uploaded to the beamformer to fulfill Equation (3).

The data of each beam is accumulated with a preset integration time, the minimum is 0.1 ms. The accumulated data is truncated into 8 bit integers, and sent to GPU server in UDP packets through the 10GbE NICs. The raw data flow is also monitored for anomalous values, which will raise alerts for possible system problems.

2.4. Data Real-Time Processing

The stream data of the digitally formed beams are further processed on the two GPU servers, each equipped with an Intel Xeon E5-2690 V4 CPU, 128 GB RAM and two NVIDIA GTX 1080Ti GPUs. The pipeline consists of five parts. The data packets sent by the FPGAs are received by a running High Availability Shared Pipeline Engine (HASHPIPE)¹⁴ program, which assembles the packets to a form appropriate for processing on the GPU. De-dispersion and FRB candidates searching are then done on the GPU. In principle, the incoherent de-dispersion and FRB searching can be done by any of the algorithm and software available, and we are working on

¹² https://casper.astro.berkeley.edu/wiki/Hardware

 $^{^{13}}$ Earlier we have also tested a system with ROACH boards (Niu et al. 2019), where a detailed description of the design of the hardware architecture and data flow was given. Although in this implementation a different board is used, the data flow structure is very similar, and we refer the reader to that paper for more details.

¹⁴ https://casper.astro.berkeley.edu/wiki/HASHPIPE



Trunc Packetizer

Figure 4. The workflow of the Tianlai dish array beamformer.



Figure 5. The HASHPIPE pipeline thread.

the development of an algorithm based on the Hough transform (Zuo & Chen 2020), but at present it is done by the software HEIMDALL¹⁵ (Barsdell et al. 2012). The candidates found by HEIMDALL are screened to remove obvious RFIs, the program for this is described later. The remaining candidates are classified as possible FRB candidates or RFIs by the deep learning program FETCH¹⁶ (Agarwal et al. 2020). In the end, the FRB candidates are plotted and their information will be saved to enable human inspection. The data processing is done with multiple threads running on the server, and flag files are used for communication between the threads.

This pipeline is capable of processing beamformed data automatically and in quasi-real time. Each GPU server is connected to one FPGA beamformer with four Gbps (Gigabit SFP+) cables, and each cable transports data from four beams.

We create four HASHPIPE threads, each processing data from one beam. A HASHPIPE thread has two data buffers and three child processes, as Figure 5 shows. Each data buffer contains three looped sub-buffers. Input thread reads data packets from network sockets and writes them into one of the sub-buffers. When a sub-buffer is full, input thread will write the data into another sub-buffer and calc thread will start processing the filled sub-buffer. Calc thread re-arranges the data structure and writes them into output data buffer so that output thread is able to write the data in the Filterbank format files for further process by HEIMDALL. To improve the pipeline's I/O performance, files are placed in RAM instead of disk. Each Filterbank file contains 201.3 seconds of data. In other words, we need to complete processing existing data within 201.3 seconds before new data arrive in order to build a quasireal-time processing pipeline. The total data stream is 159.1 MB/s.

¹⁵ https://sourceforge.net/projects/heimdall-astro/

¹⁶ https://github.com/devanshkv/fetch



Figure 6. Result of ϕ_{ab} for two sequential calibrations. The solid line shows the first calibration and dashed line shows the second. Each figure show four feed channels, malfunctioning ones are not shown. The reference channels 5X and 5Y are not shown.

3. Experiment

3.1. Calibration

The complex gain g_a needs to be determined by a calibration procedure. One possibility is to use the original correlator to do this calibration (Zuo et al. 2019), but the separate circuit for the original correlator and the FRB beam former may introduce some differences. We have therefore adopted a self-calibration approach based entirely on the FRB beam former.

The absolute value of each complex gain is first adjusted so that the amplitude of the readout for each feed is approximately equal (within 1 dB). The instrument phase is then calibrated by observing the transit of a bright point source, during which we make special beams used for calibration. In each such special beam, all weights are set to be zero except for a reference antenna and the antennas to be calibrated. In the ground coordinate system, the antennas are fixed while the source moves with the rotation of celestial sphere, so the output of this beam for a pair of antennas a, b is

$$I_{ab}(t) = |g_a|^2 |A_a(\mathbf{n})|^2 + |g_b|^2 |A_b(\mathbf{n})|^2 + |g_a||g_b|[A_a(\mathbf{n})A_b^*(\mathbf{n})e^{-j[2\pi(\mathbf{n}-\mathbf{k})\cdot\mathbf{u}_{ab}+(\phi_a-\phi_b)]} + \text{c.c.}]$$
(9)

where c.c. denotes complex conjugate, $\boldsymbol{u}_{ab} = \boldsymbol{u}_b - \boldsymbol{u}_a$, and $\boldsymbol{n}(t)$ is the position of point source at time *t* which is known (we have omitted to write out *t*), and \boldsymbol{k} is the direction of the primary beam center. Denote $\varphi_{ab}(t) = 2\pi\boldsymbol{n}(t) \cdot \boldsymbol{u}_{ab}$, we model the primary beam *A* as a sinc function, fitting Equation (9) by:

$$\operatorname{Re} I_{ab}(t) = A^{2} [\cos(\varphi_{ab}(t) - \varphi_{ab}(t_{\operatorname{transit}}) + \phi_{ab}) + C^{2}] \\ \times \operatorname{sinc}^{2} \left(\frac{\varphi_{ab}(t) - \varphi_{ab}(t_{\operatorname{transit}})}{2\sigma} \right),$$
(10)



Figure 7. The output amplitude of the 16 beams as a function of time. We set the beams along the track of Cas A with a spacing of 0.5.



Figure 8. Beam amplitude vs. time, where there are 14 beams synthesized from signal from 1 to 14 antennas, respectively.

where t_{transit} is the time for the source transit (passing the beam center). The instrument phase difference $\phi_{ab} = \phi_a - \phi_b$ is then determined, assuming it to be constant during the calibration process.

The beamformer allows us to simultaneously form up to 16 beams for each polarization. From each of these the instrument phase difference of that pair can be solved by the transit observation. If the set of pairs include all 16 antennas, then in principle the instrument phase of all 16 antennas can be solved, up to an overall arbitrary phase offset, which can be set to zero for the reference antenna. However, we need to use baselines which have large variation of geometric phase during the transit to obtain good solution for ϕ_{ab} . Due to the circular geometry of the Tianlai array, this is not possible for a single reference antenna. Instead, we choose two reference antennas during each calibration. For example, in one calibration of Cas A transit, we can choose antennas 1 and 5 which are located at opposite sides of the array as reference antennas. The remaining antennas are either paired with antenna 1 or antenna 5, whichever allows a longer projected baseline for the transit

 Table 2

 Beam Output vs. Number of Antenna in Formed Beam

Antenna Number	Prediction	Measurement	Efficiency
1	1	1	1
2	4	3.23	0.81
3	9	6.91	0.77
4	16	12.4	0.78
5	25	20.1	0.80
6	36	30.1	0.84
7	49	40.6	0.83
8	64	52.1	0.81
9	81	66.1	0.82
10	100	81.6	0.82
11	121	95.1	0.79
12	144	113	0.78
13	169	133	0.79
14	196	142	0.72

position. The instrument phases of the antennas with respect to either antenna 1 or antenna 5 are then obtained. The antennas 1 and 5 are also paired, so that the instrument phase difference between the two is obtained. Then the phase of every antenna of the array can be solved with respect to the designated antenna.

In Figure 6, we show the instrument phases with respect to antenna 5 for each feed in two sequential calibrations separated by an hour. To check the reliability of the calibration, we have used two different sets of references. In the first observation, antennas 3 and 14 are used as reference antennas, while in the second one antennas 1 and 5 are used. The phase obtained from the two calibrations agree very well, showing that our results are consistent for whichever pair of reference antennas, and the instrument remains very stable during night time.

With the instrument phase determined, we can now form beams using all antennas and make observation. We test our beam pointing accuracy by making beams along the track of a bright source, with a spacing of 0°.5. We expect to see the output amplitude of each formed beam reaches peak one by one as the order of beam number. The result is shown in Figure 7 as anticipated.

If the complex gains of all antennas are perfectly accurate, the output of all antennas will add coherently, then for *n* antennas, the output for a source at the beam center should $\propto n^2$. In reality, there will be errors in the calibration of the amplitudes and phases of the antennas, then when we form the beams, the output would be less than that of n^2 . We test the efficiency of beamforming by observing the strong radio source Cas A. We first use the transit of Cas A to calibrate the complex gains. Then, we re-point our antennas so that the Cas A can pass them again. During this second transit, we form the beams as follows: all beams are pointing to the same direction, beam #1 has non-zero weight for antenna 1 only; beam #2 has non-

Table 3HEIMDALL Configuration

	-
SNR threshold	6
DM range	10-2000
SNR loss tolerance	1.5
DM trials	866
Max boxcar width	128
Max giant rate	2000 per minute

zero weight for antennas 1 and 2; beam #3 has non-zero weight for antennas 1, 2, 3; and so on. As we have two malfunctioned antennas during the test, we formed up to 14 beams, with beam #14 include all 14 working antennas. The transit curves are shown in Figure 8. The expected output, the measured output, and the beamforming efficiency (defined as measured/expect) are given in Table 2. The typical efficiency achieved is about 0.8.

3.2. De-dispersion

We use HEIMDALL, a GPU accelerated radio transient detection pipeline, to search for single pulses in our data. For m DM trials and n boxcar width trials, HEIMDALL creates $m \times n$ de-dispersed and downsampled time series to cover parameter space. For each time series, the median value of the data is subtracted and the data is divided by the standard deviation so as to convert amplitude data into signal-to-noise ratio (SNR) data. Samples in the SNR data that exceed a preset threshold, usually 6, are marked as "giants". Giants from different beams are checked for coincidence, those with the same time and DM are joined as one candidate. Finally, HEIMDALL outputs all candidates found in a candidate file.

The HEIMDALL parameters we set are listed in Table 3. For candidates signal-to-noise ratio threshold, we use 6, the default setting. DM search range is set to be 10-2000. Most RFIs can be excluded by the lower DM limit, as they do not exhibit dispersion. The upper limit is set to limit the required computation. The Tianlai dish array is a very small one, with quite limited sensitivity, so it is not very likely to detect very distant FRBs with large dispersion. The SNR loss tolerance is the maximum SNR loss allowed between DM trails. Under our setting, HEIMDALL has a total of 874 DM trials for each data set. The maximum boxcar width sets the upper limit of pulse width that HEIMDALL searches for. It is set to be 128 samples, which means we search pulse width up to 25.2 milliseconds. Because strong RFIs may create huge amount of candidates and lead to pipeline failure, a maximum giant rate is set to restrict the maximum candidates per minute that HEIMDALL reports. If the number of candidates exceeds the threshold, HEIMDALL will abort further searching and report the found candidates.



Figure 9. Overall workflow of our pipeline.



Figure 10. The FETCH candidate plot of a pulse from PSR B0329+54. Left: time-DM plot. Right: time-frequency (waterfall) plot. The band at 760 MHz is RFI.



Figure 11. The dynamic spectrum of the pulse in Figure 10. Left: before de-dispersion. Right: after de-dispersion.

3.3. Candidates Filtering and Classification

With these settings, HEIMDALL reports thousands of candidates with S/N > 6 per minute. We further filter these candidates out to remove obvious RFIs. The pipeline is shown in Figure 9.

Candidates with total S/N < 7 are removed, and those with DM $< 20 \text{ pc cm}^{-2}$ are considered to be RFI. The DM criteria here are set at a higher value than the one set for HEIMDALL, in order to remove very strong and broad RFIs, which may still be strong enough to become candidates when the data are dedispersed to DM = 10.

As we form our beams in different directions while searching for FRBs, candidates found in all beams at the same time are most likely an artificial radio frequency interference, which is either in the near field of the array, or too strong. Such events are filtered out. If all candidates within a time frame are rejected by the filter program, the data files in the RAM are removed. Data files with candidates passing through the filter are saved to the disk.

We next use the Fast Extragalactic Transient Candidate Hunter (FETCH) (Agarwal et al. 2020), a deep-learning based classifier, to classify the remaining candidates. FETCH takes the SNR of the candidate, the starting time of the burst, the DM value and the box-car filter width reported by the HEIMDALL program as its input. It plots the output amplitude distribution in the time-DM and time-frequency figure for each candidate. An example is shown in Figure 10 for the pulsar test observation (described below). A pre-trained deep learning model takes these two figures as input and output the possibility of the candidates being a genuine astronomical source. Candidates with possibility greater than 50% will be regarded as plausible FRBs. Data files with no plausible candidates are removed.

3.4. Pulse Searching Test

We tested our system with known pulsars. The pointing of our beams is put to test by observing the brighter pulsar PSR B0329+54. We evenly spaced our 16 beams along the pulsar's transit track with a separation of 0°.5, and then carried out an observation. The data is run through the pipeline, and the pulses from the pulsar are successfully detected. In Figure 11, we show the dynamic spectrum of a pulse before (left) and after de-dispersion.



Figure 12. De-dispersed time series observing PSR B2021+51 (top) and B0329+54 (bottom).

We further validate the system by observing PSR B0329+54 and PSR B2021+51. We succeed in detecting pulses from both pulsars, as shown in Figure 12. As our beams overlapped with adjacent ones, pulsars are visible in two or three beams simultaneously during observation.

3.5. Sensitivity

We can estimate the sensitivity of system by using the data in Section 3.1 and Equation (11),

Sensitivity =
$$\frac{S_{M1}}{I_{\text{peak}} - I_{\text{noise}}} * \Delta t * 3\sigma_{\text{noise}}$$
 (11)

our 3σ fluence sensitivity at the center of primary beam is estimated to be 7.6 Jy ms.

We arrange our 16 digital beams in the following way: one is centered at the NCP during our regular survey, and the other 15 beams as two concentric rings around the center beam, one configuration (the optimized one) is shown in Figure 13. To maximize our FRB detection rate, we optimized three parameters: the number of beams in the inner circle, and the diameters of the inner and outer circles. We create a function that takes the above parameters as input and returns the estimated overall FRB detection rate.

To obtain an estimated FRB detection rate for a given configuration, first we calculate the position of the beam centers, then obtain each beam's profile by shifting the digital beam profile to the given beam center, and multiplying it with the primary beam profile. For a given pixel, we estimate its effective detecting beam profile by using the beam with the best gain on that pixel, usually this is the beam with the nearest beam center. The FRB rate of every pixel is then calculated based on this overall beam profile.

We take the all sky FRB rate as 820 per day above 5 Jyms at 400–800 MHz and fluence index $\alpha = -1.4$, as estimated from CHIME observations (CHIME/FRB Collaboration et al. 2021). The overall FRB detection rate is obtained by summing over all pixels. Last, we find the maximum value of this function by varying the parameters. The optimal configuration with the above general lay out is to have six beams in the inner circle, and nine beams in the outer circle, with diameters of



Figure 13. Overall beam profile. The red circle indicates the -3 dB area of the primary beam.



Figure 14. The expected differential detection rate of 10σ FRBs as a function of fluence. Each bar indicates FRB rate within 5 Jyms. The FRBs with fluence greater than 300 Jyms are not shown in the figure but has a total rate of ~0.05 month⁻¹.

1°.46 and 2°.82 respectively. The corresponding FRB rate is ~ 1.2 FRBs per month for 3σ bursts and ~ 0.27 FRBs per month for 10σ bursts, and the expected distribution of 10σ FRB fluence is shown in Figure 14.

Note that this is a rough estimate based on simplified treatment, because in reality, the signal-to-noise ratio depends not only on the fluence of the burst, but also on the shape of the signal, the distribution of the noise over the frequency band, and the RFI. We have also ignored the additional information provided by nearby beams, which could be important for the pixels near the border of adjacent beams.

4. Conclusion

We have designed and installed the FRB-searching backend for the Tianlai Dish Pathfinder Array, using a combination of FPGA beamformer, and an off-shelf GPU server for de-dispersion. The system can form 16 beams for each polarization. The system has been validated with test observation of pulsars. The Tianlai Dish Pathfinder Array will conduct a deep survey of the north celestial polar region by continuously pointing its telescopes to the NCP region (Perdereau et al. 2022), and the FRB search can be conducted concurrently. The data processing pipeline can make real-time processing of the data, and automatically raise alerts for possible follow up observation by other telescopes. The estimated 10σ sensitivity at the beam center is 25.3 Jy ms. For an FRB distribution estimated by CHIME/FRB Collaboration et al. (2021), a simple estimate gives an expected FRB detection rate of ~0.27 FRBs per month.

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